

# **FTCS Solution to the Heat Equation**

**ME 448/548 Notes**

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## Overview

1. Use the forward finite difference approximation to  $\partial u / \partial t$ .

$$\frac{\partial u}{\partial t} \approx \frac{u_i^{k+1} - u_i^k}{\Delta t}$$

2. Use the central difference approximation to  $\partial^2 u / \partial x^2$  at time  $t_k$ .

$$\left. \frac{\partial^2 u}{\partial x^2} \right|_{x_i} = \frac{u_{i-1}^k - 2u_i^k + u_{i+1}^k}{\Delta x^2}$$

3. Solve for  $u_i^{k+1}$ . The computational formula is *explicit*: each value of  $u_i^{k+1}$  can be updated independently.
4. FTCS is easy to implement, but is only conditionally stable
5. Truncation errors are  $\mathcal{O}((\Delta x)^2)$  and  $\mathcal{O}(\Delta t)$ .

## Finite Difference Operators

Choose the *forward difference* to evaluate the time derivative at  $t = t_k$ .

$$\left. \frac{\partial u}{\partial t} \right|_{t_k, x_i} = \frac{u_i^{k+1} - u_i^k}{\Delta t} + \mathcal{O}(\Delta t) \quad (1)$$

Approximate the spatial derivative with the central difference operator and take all nodal values at time  $t_k$ .

$$\left. \frac{\partial^2 u}{\partial x^2} \right|_{x_i} = \frac{u_{i-1}^k - 2u_i^k + u_{i+1}^k}{\Delta x^2} + \mathcal{O}(\Delta x^2) \quad (2)$$

## FTCS Approximation to the Heat Equation

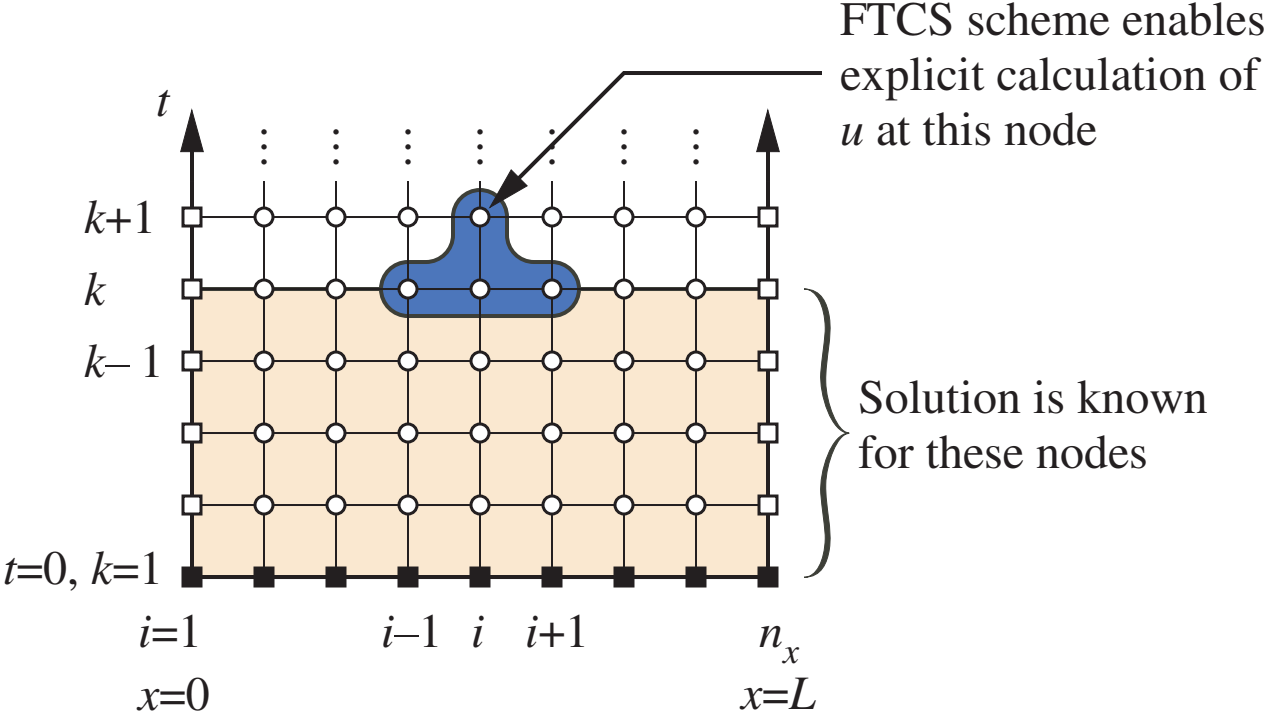
Substitute Equation (1) and Equation (2) into the heat equation

$$\frac{u_i^{k+1} - u_i^k}{\Delta t} = \alpha \frac{u_{i-1}^k - 2u_i^k + u_{i+1}^k}{\Delta x^2} + \mathcal{O}(\Delta t) + \mathcal{O}(\Delta x^2) \quad (3)$$

Drop truncation error terms to get

$$\frac{u_i^{k+1} - u_i^k}{\Delta t} = \alpha \frac{u_{i-1}^k - 2u_i^k + u_{i+1}^k}{\Delta x^2} \quad (4)$$

# FTCS Computational Molecule



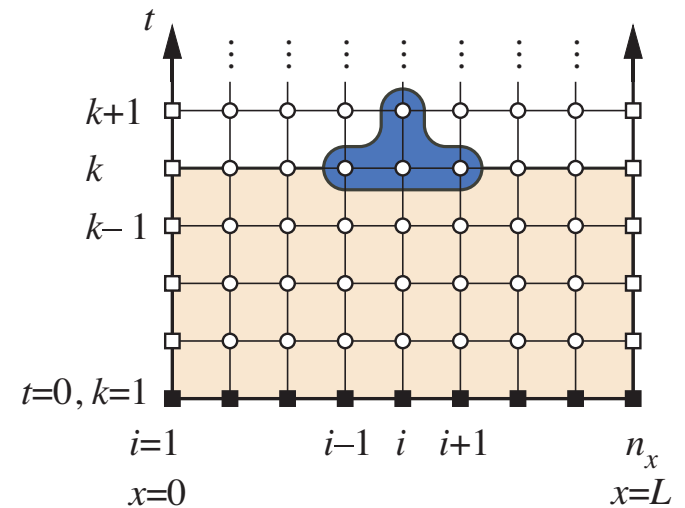
## FTCS Approximation to the Heat Equation

Solve Equation (4) for  $u_i^{k+1}$

$$u_i^{k+1} = ru_{i+1}^k + (1 - 2r)u_i^k + ru_{i-1}^k \quad (5)$$

where  $r = \alpha\Delta t/\Delta x^2$ .

FTCS is an *explicit* scheme because it provides a simple formula to update  $u_i^{k+1}$  independently of the other nodal values at  $t_{k+1}$ .



## demoFTCS Code

```
function errout = demoFTCS(nx,nt)

% ... Comments and processing of optional inputs skipped

% --- Assign physical and mesh parameters
alfa = 0.1; L = 1; tmax = 2; % Diffusion coefficient, domain length and max time
dx = L/(nx-1); dt = tmax/(nt-1);
r = alfa*dt/dx^2; r2 = 1 - 2*r;

% --- Assign IC and BC. u is initialized to a vector that includes BC
x = linspace(0,L,nx)'; u = sin(pi*x/L);

% --- Loop over time steps
for k=2:nt
    uold = u; % prepare for next step
    for i=2:nx-1
        u(i) = r*uold(i-1) + r2*uold(i) + r*uold(i+1);
    end
end
```

Remember that the formula for updating  $u_i^{k+1}$  is

$$u_i^{k+1} = ru_{i+1}^k + (1 - 2r)u_i^k + ru_{i-1}^k$$

## Alternative formulation to the FTCS Algorithm

Equation (5) can be expressed as a matrix multiplication.

$$u^{(k+1)} = Au^{(k)} \quad (6)$$

where  $u^{(k+1)}$  is the vector of  $u$  values at time step  $k + 1$ ,  $u^{(k)}$  is the vector of  $u$  values at time step  $k$ , and  $A$  is the *tridiagonal matrix*

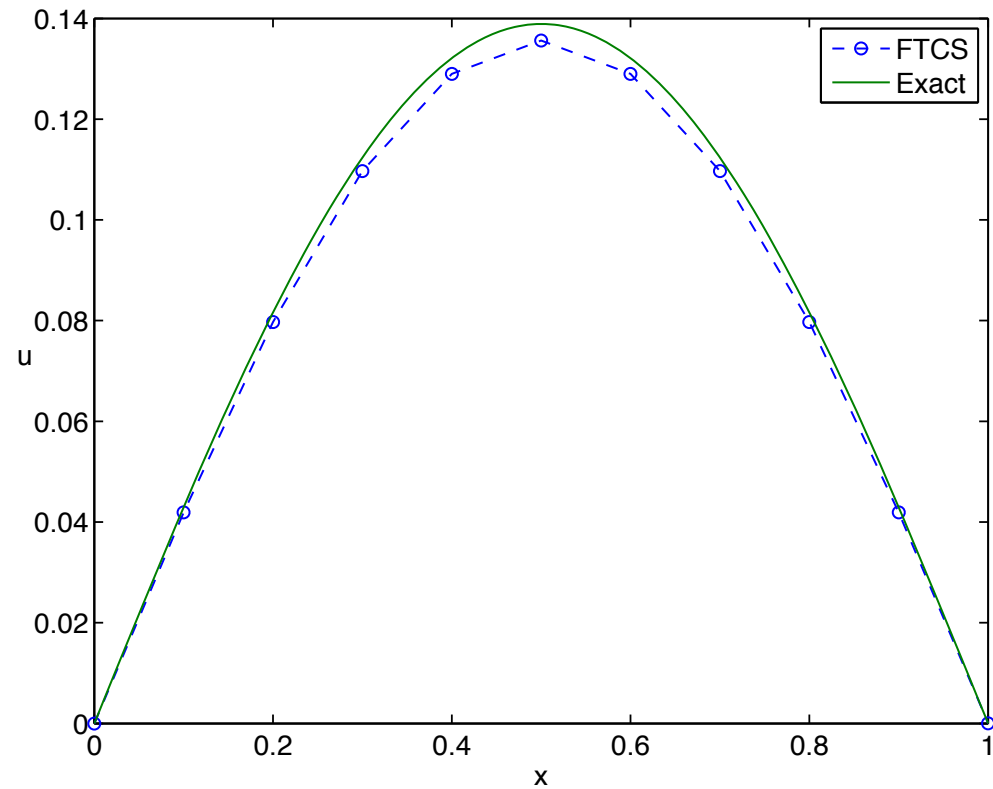
$$A = \begin{bmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ r & (1 - 2r) & r & 0 & 0 & 0 \\ 0 & r & (1 - 2r) & r & 0 & 0 \\ 0 & 0 & \dots & \dots & \dots & 0 \\ 0 & 0 & 0 & r & (1 - 2r) & r \\ 0 & 0 & 0 & 0 & 0 & 1 \end{bmatrix}. \quad (7)$$

The first and last rows of  $A$  are set to enforce the Dirichlet boundary conditions at  $x = 0$  and  $x = L$ .



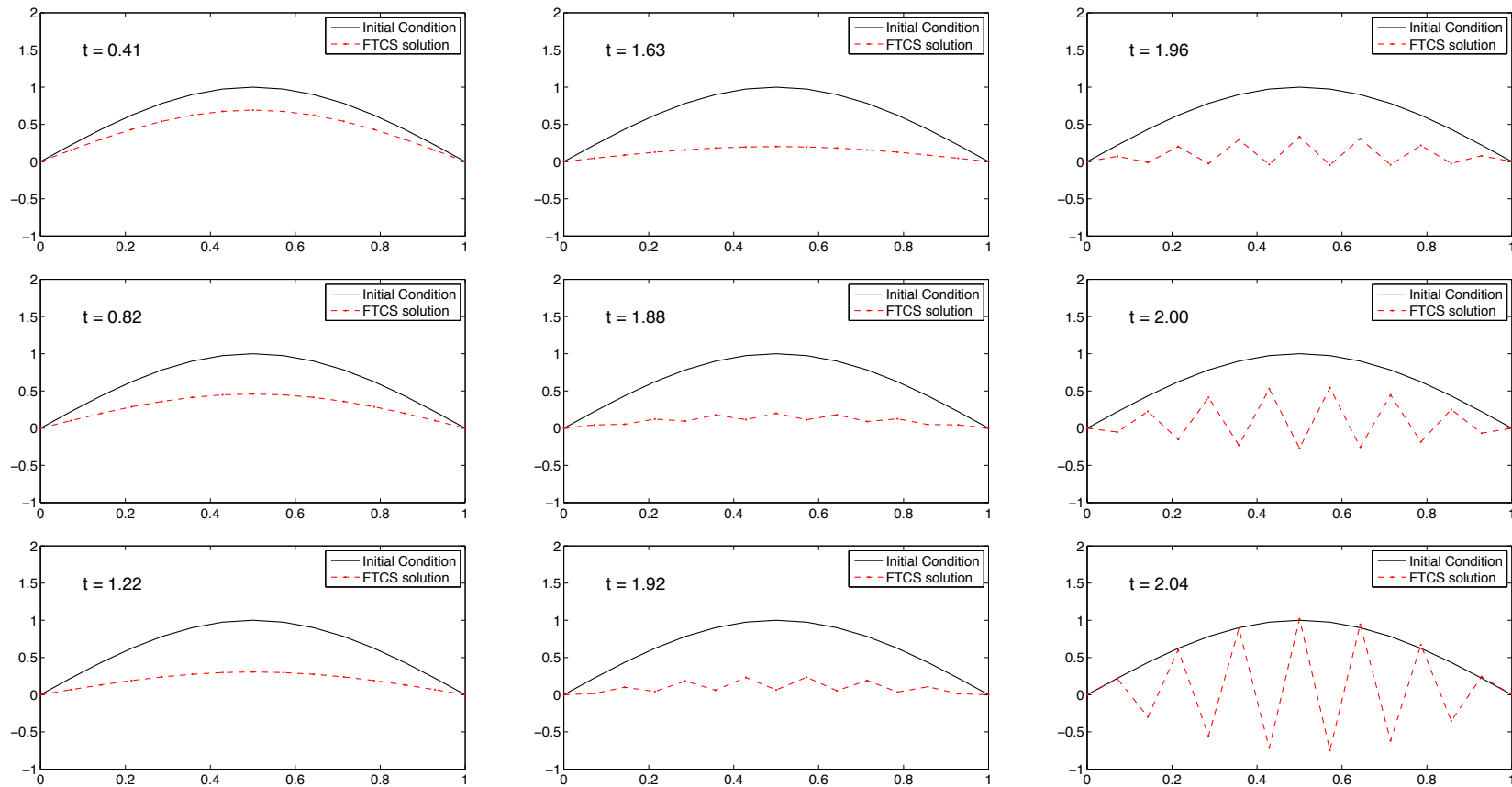
## Run the demoFTCS code

```
>> demoFTCS  
Error in FTCS solution = 0.002221
```



# Unstable FTCS Solution

movieFTCS(15):



## Conditionally Stable FTCS Scheme

Observations:

- The FTCS solution appears to be stable at first
- Toward the end of the simulation time, oscillations grow exponentially
- Instability is not caused by truncation error
- Instability is fed by round-off errors, but not directly caused by round-off

We'll use a simplified form of Fourier Stability Analysis

- Suppose that the initial condition is a small sine wave
- The correct solution is a decay of the sine wave
- Under what condition does the solution grow instead of decay?

*Stability or instability is an intrinsic property of the scheme*

FTCS is *conditionally stable* for solutions to the heat equation.

## Stability Analysis

Suppose the initial condition looks like this

$$u_0(x) = \sigma \cos\left(\frac{\pi x}{\Delta x}\right) = \sigma(-1)^{i-1}$$

where  $\sigma$  is a suitably small value

The solution at time step  $k = 2$  is

$$\begin{aligned} u_i^{(2)} &= r(u_0(x_{i+1}) + u_0(x_{i-1})) + (1 - 2r)u_0(x_i) \\ &= r\sigma((-1)^i + (-1)^{i-2}) + (1 - 2r)\sigma(-1)^{i-1} \\ &= -r\sigma((-1)^{i+1} + (-1)^{i-1}) + (1 - 2r)\sigma(-1)^{i-1} \end{aligned}$$

## Stability Analysis

Since  $(-1)^{i+1} = (-1)^{i-1}$  for any  $i$ , the solution at  $k = 2$  can be further simplified

$$\begin{aligned}u_i^{(2)} &= -2r\sigma(-1)^{i-1} + (1 - 2r)\sigma(-1)^{i-1} \\&= (1 - 4r)\sigma(-1)^{i-1} \\&= (1 - 4r)u_0(x_i)\end{aligned}$$

The pattern is

$$\begin{aligned}u_i^{(2)} &= (1 - 4r)u_0(x_i) \\u_i^{(3)} &= (1 - 4r)u_i^{(2)}(x_i) = (1 - 4r)^2u_0(x_i) \\u_i^{(4)} &= (1 - 4r)u_i^{(3)}(x_i) = (1 - 4r)^3u_0(x_i)\end{aligned}$$

where  $u_i^{(2)}$  is the numerical solution at  $x_i$  and  $k = 2$ ;  
and where  $(1 - 4r)^2$  is the square of  $(1 - 4r)$

## Stability Analysis

The general pattern is

$$u_i^k = (1 - 4r)^{k-1} u_0(x_i).$$

where  $(1 - 4r)^{k-1}$  is  $(1 - 4r)$  raised to the  $k - 1$  power.

Therefore, the solution grows when  $|1 - 4r| > 1$ .

Therefore, stability requires  $1 - 4r < 1$  and  $-(1 - 4r) < 1$ .

$$1 - 4r < 1 \implies -4r < 2 \implies r > -\frac{1}{2} \text{ true for any } r > 0.$$

The second case is

$$-(1 - 4r) < 1 \implies 4r < 2 \implies r < \frac{1}{2}$$

## Stability Analysis

The quick and dirty stability analysis shows that the FTCS scheme is stable only if

$$r = \frac{\alpha \Delta t}{\Delta x^2} < \frac{1}{2}. \quad (8)$$

Recall:  $\alpha$  is a parameter of the physical problem.

We must choose  $\Delta t$  and  $\Delta x$  so that Equation (8) is satisfied.

## Working with the FTCS Stability Criterion

To increase accuracy, we want to decrease both  $\Delta x$  and  $\Delta t$ .

For a given  $\Delta x$ , the stability limit for FTCS imposes an upper limit on  $\Delta t$

$$\Delta t < \frac{\Delta x^2}{2\alpha} \quad (9)$$

Choosing  $\Delta t$  and  $\Delta x$  so that  $r < \frac{1}{2}$  does not guarantee an accurate numerical solution.

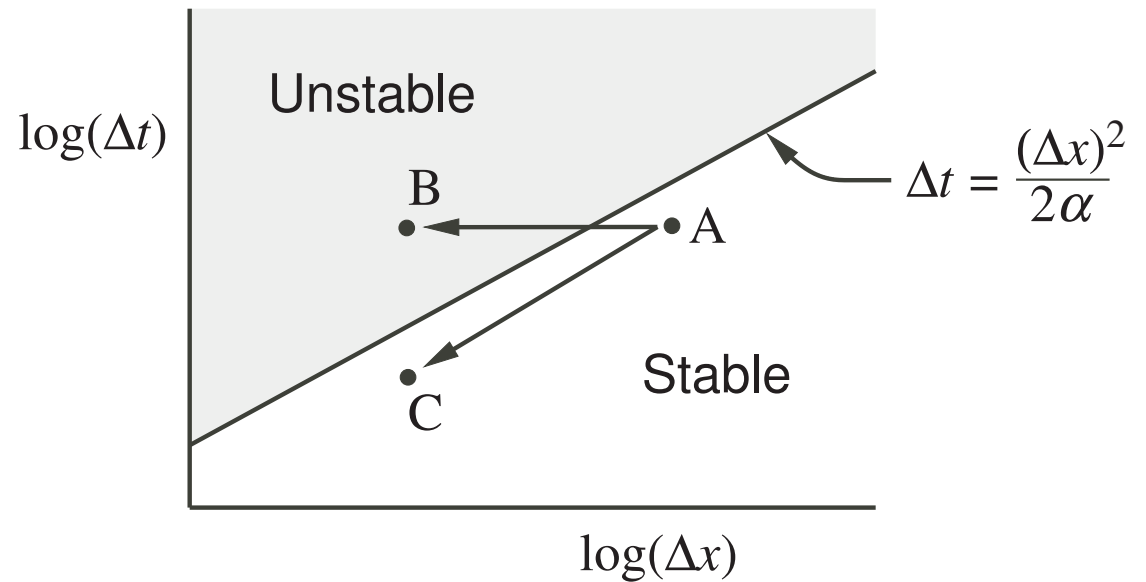
$r < \frac{1}{2}$  only guarantees that the FTCS solution will not blow up.



## Working with the FTCS Stability Criterion

When reducing  $\Delta x$  and  $\Delta t$  to improve accuracy, follow a path like  $A \rightarrow C$  not  $A \rightarrow B$ .

Note that the axes have logarithmic scales.



## Measuring the FTCS Truncation Error

The error reported by demoFTCS is defined as

$$E(n_x, n_t) = \frac{1}{\sqrt{n_x}} \|u_i^k - u(x_i, t_k)\|_2 \quad (10)$$

The factor of  $1/\sqrt{n_x}$  converts  $\|u_i^k - u(x_i, t_k)\|_2$  to an average error.

## Measuring the FTCS Truncation Error

Designate the local error at  $x = x_i$  and  $t = t_k$  as

$$e_i^k = u_i^k - u(x_i, t_k). \quad (11)$$

Define  $\bar{e}_k$  as an *RMS average error per node* at time step  $t_k$

$$\bar{e}^k \equiv \left[ \frac{1}{n_x} \sum_{i=1}^{n_x} (e_i^k)^2 \right]^{1/2} \quad (12)$$

Algebraic substitution shows that  $E(n_x, n_t) = \bar{e}^k$ .

If the solution to the heat equation is smooth, then  $E(n_x, n_t) = \mathcal{O}(\Delta t) + \mathcal{O}(\Delta x^2)$

## Measuring the truncation error

Systematically reduce  $\Delta x$  and  $\Delta t$  to determine whether truncation error prediction is realized by the code.

We know the exact solution, so we use Equation(10).

*This is a test* of whether the code correctly implements FTCS.

Recall that for FTCS

$$E(n_x, n_t) = \mathcal{O}(\Delta t) + \mathcal{O}(\Delta x^2)$$

but stability requires  $\Delta t = C \Delta x^2$ , where  $C$  is a constant.

Therefore, a stable solution to FTCS should demonstrate

$$E(n_x, n_t) = \mathcal{O}(\Delta x^2)$$

## Measuring the truncation error

Suppose we don't know the exponent of the truncation error for our code.

In other words, instead of

$$E(n_x, n_t) = \mathcal{O}(\Delta x^2)$$

we have

$$E(n_x, n_t) = \mathcal{O}(\Delta x^p)$$

where  $p$  is unknown

Since  $\Delta x = \frac{L}{n_x - 1}$  we have

$$E(n_x, n_t) = \mathcal{O}(\Delta x^p) = \mathcal{O}\left(\frac{L^p}{(n_x - 1)^p}\right)$$

## Measuring the truncation error

Simplify the preceding expression

$$\begin{aligned} E(n_x, n_t) &= \mathcal{O}(\Delta x^p) \\ &= \mathcal{O}\left(\frac{L^p}{(n_x - 1)^p}\right) \\ &\sim \mathcal{O}\left(\frac{1}{(n_x - 1)^p}\right) \quad \text{ignore multiplicative constants} \\ &\sim \mathcal{O}\left(\frac{1}{n_x^p}\right) \quad \text{leading powers dominate} \end{aligned}$$

Therefore, we expect

$$E(n_x, n_t) = \mathcal{O}\left(\frac{1}{n_x^p}\right)$$

## Measuring the truncation error

Measure  $E(n_x, n_t)$  on two difference meshes,  $n_{x,1}$  and  $n_{x,2}$ .

With

$$E(n_x, n_t) = \mathcal{O}\left(\frac{1}{n_x^p}\right)$$

we form the ratio

$$\frac{E(n_{x,2}, n_t)}{E(n_{x,1}, n_t)} = \frac{n_{x,1}^p}{n_{x,2}^p} = \left(\frac{n_{x,1}}{n_{x,2}}\right)^p$$

Solve for  $p$

$$p = \frac{\log(E(n_{x,2}, n_t)/E(n_{x,1}, n_t))}{\log(n_{x,1}/n_{x,2})}.$$

## convFTCS Code

```
function convFTCS
% convFTCS Convergence of FTCS on a series of finer spatial meshes

% --- Set constants to be consistent with demoFTCS
alfa = 0.1; L = 1; tmax=2; rsafe = 0.49999; % stable r<0.5

% --- Specify nx and compute nt consistent with stability limit
nx = [8 16 32 64 128 256]; dx = L./(nx-1);
nt = ceil( 1 + alfa*tmax*((nx-1).^2)/(rsafe*L^2) );

% --- Loop over mesh sizes, store error and compute order of scheme
fprintf('\n nx nt error E(j)/E(j-1) p\n');
er = NaN; p = 0;
for j=1:length(nx);
    e(j) = demoFTCS(nx(j),nt(j));
    if j>1
        er = e(j)/e(j-1); p = log(er)/log(nx(j-1)/nx(j));
    end
    fprintf(' %5d %5d %11.3e %8.4f %8.4f\n',nx(j),nt(j),e(j),er,p);
end

% -- plotting code skipped ...
```

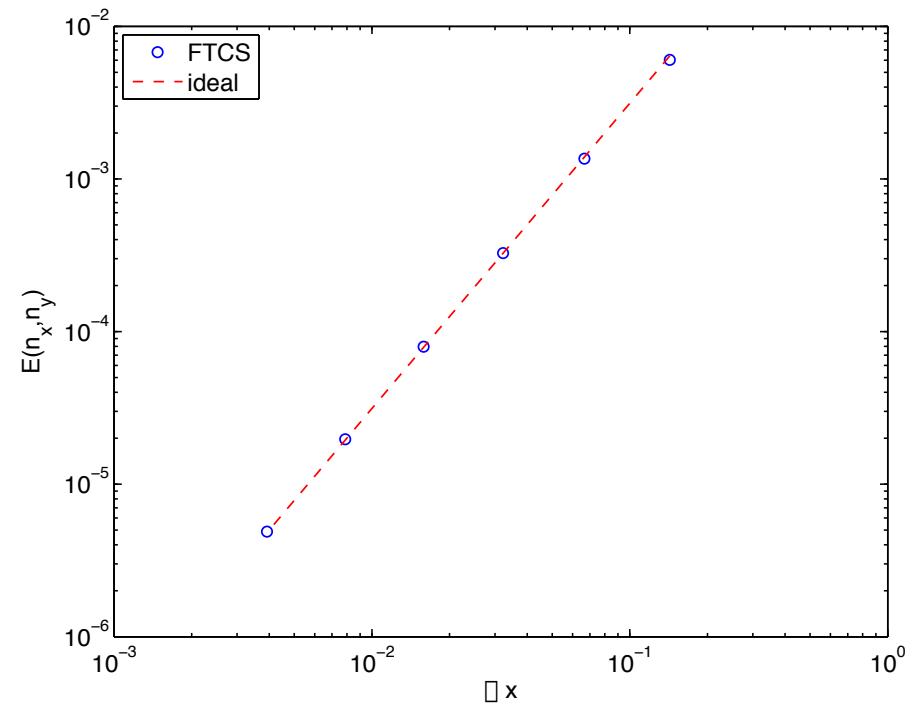


## Running the convFTCS Code

```
>> convFTCS
```

| nx  | nt    | error     | E(j)/E(j-1) | p      |
|-----|-------|-----------|-------------|--------|
| 8   | 21    | 6.028e-03 | NaN         | 0.0000 |
| 16  | 92    | 1.356e-03 | 0.2249      | 2.1524 |
| 32  | 386   | 3.262e-04 | 0.2406      | 2.0553 |
| 64  | 1589  | 7.972e-05 | 0.2444      | 2.0329 |
| 128 | 6453  | 1.970e-05 | 0.2471      | 2.0170 |
| 256 | 26012 | 4.895e-06 | 0.2485      | 2.0085 |

The last column of text output shows that demoFTCS has the right truncation error behavior.



## Summary for the FTCS Scheme

- FTCS is easy to implement. The update formula is

$$u_i^{k+1} = ru_{i+1}^k + (1 - 2r)u_i^k + ru_{i-1}^k$$

- The FTCS scheme is conditionally stable when

$$r = \frac{\alpha \Delta t}{\Delta x^2} < 1/2$$

In two spatial dimensions with  $\Delta y = \Delta x$ , the stability condition is  $r < 1/4$ . In 3D with  $\Delta y = \Delta z = \Delta x$ , the stability condition is  $r < 1/8$ .

- There are *much better* schemes for solving the heat equation.
- FTCS is a toy used to introduce the numerical solution of PDEs