

Crank Nicolson Solution to the Heat Equation

ME 448/548 Notes

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Overview

1. Use finite approximations to $\partial u / \partial t$ and $\partial^2 u / \partial x^2$: same components used in FTCS and BTCS.
2. Average contribution from t_k and t_{k+1} in the approximation of $\partial^2 u / \partial x^2$.
3. Computational formula is (still) *implicit*: all u_i^{k+1} must be solved simultaneously. Information from u_{i-1}^k , u_i^k , and u_{i+1}^k are used at each time step in the computation of u_i^{k+1} .
4. Solution is not significantly more complex than BTCS.
5. Like BTCS, the Crank-Nicolson method is unconditionally stable for the heat equation.
6. Truncation error is $\mathcal{O}(\Delta t^2) + \mathcal{O}(\Delta x^2)$. Time accuracy is better than BTCS or FTCS.

The θ Method

Evaluate the diffusion operator $\partial^2 u / \partial x^2$ at both time steps t_{k+1} and time step t_k , and use a weighted average

$$\frac{u_i^{k+1} - u_i^k}{\Delta t} = \alpha\theta \left[\frac{u_{i-1}^{k+1} - 2u_i^{k+1} + u_{i+1}^{k+1}}{\Delta x^2} \right] + \alpha(1 - \theta) \left[\frac{u_{i-1}^k - 2u_i^k + u_{i+1}^k}{\Delta x^2} \right] \quad (1)$$

where

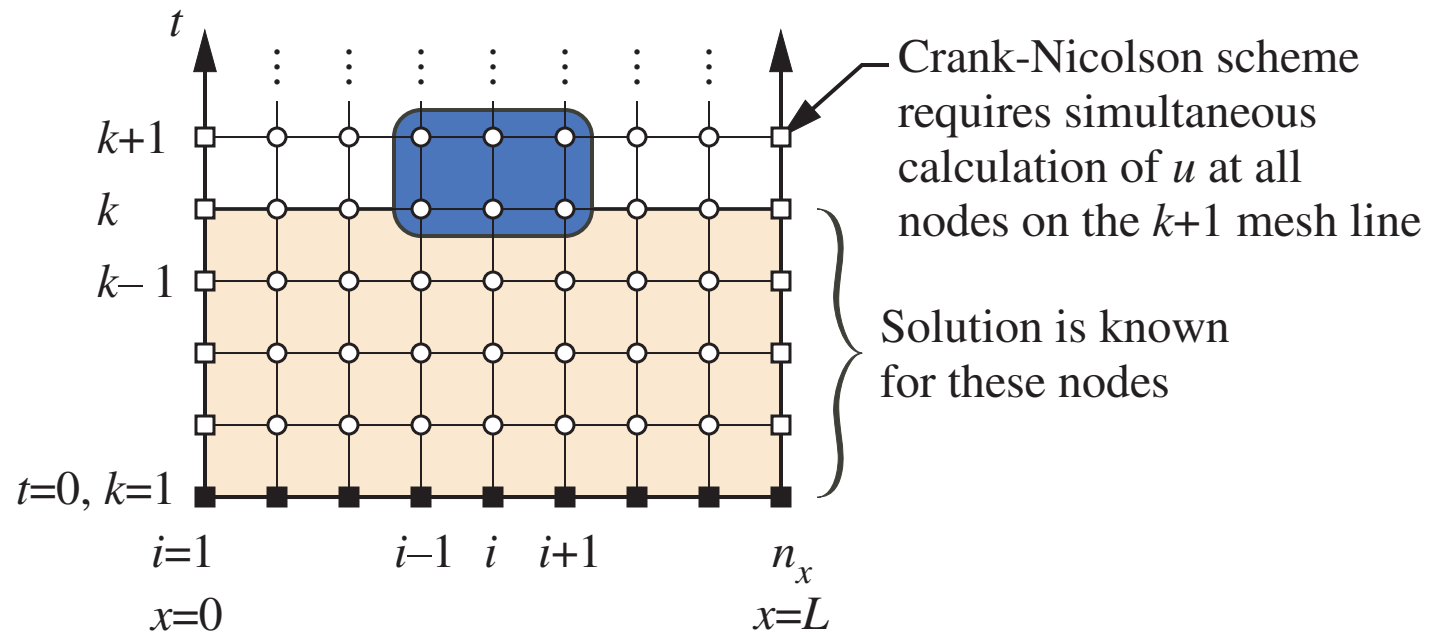
$$0 \leq \theta \leq 1$$

$$\theta = 0 \quad \iff \quad \text{FTCS}$$

$$\theta = 1 \quad \iff \quad \text{BTCS}$$

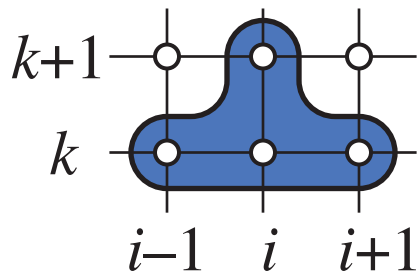
$$\theta = \frac{1}{2} \quad \iff \quad \text{Crank-Nicolson}$$

Crank-Nicolson Computational Molecule

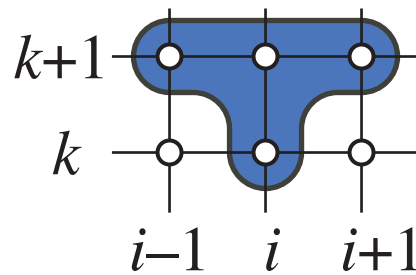


Compare Computational Molecules

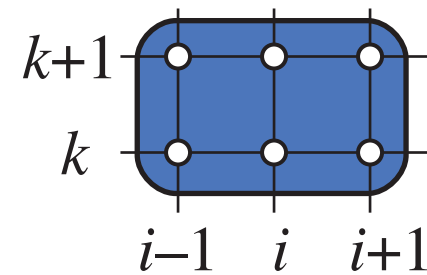
FTCS



BTCS



Crank-Nicolson



Crank Nicolson Approximation to the Heat Equation

Set $\theta = \frac{1}{2}$ in the formulation of the θ method.

$$-\frac{\alpha}{2\Delta x^2}u_{i-1}^{k+1} + \left(\frac{1}{\Delta t} + \frac{\alpha}{\Delta x^2}\right)u_i^{k+1} - \frac{\alpha}{2\Delta x^2}u_{i+1}^{k+1} =$$
$$\frac{\alpha}{2\Delta x^2}u_{i-1}^k + \left(\frac{1}{\Delta t} - \frac{\alpha}{\Delta x^2}\right)u_i^k + \frac{\alpha}{2\Delta x^2}u_{i+1}^k \quad (2)$$

Crank-Nicolson System of Equations

The system of equations has the same structure as BTCS

$$\begin{bmatrix} a_1 & b_1 & 0 & 0 & 0 & 0 \\ c_2 & a_2 & b_2 & 0 & 0 & 0 \\ 0 & c_3 & a_3 & b_3 & 0 & 0 \\ 0 & 0 & \cdots & \cdots & \cdots & 0 \\ 0 & 0 & 0 & c_{n_x-1} & a_{n_x-1} & b_{n_x-1} \\ 0 & 0 & 0 & 0 & c_{n_x} & a_{n_x} \end{bmatrix} \begin{bmatrix} u_1^{k+1} \\ u_2^{k+1} \\ u_3^{k+1} \\ \vdots \\ u_{n_x-1}^{k+1} \\ u_{n_x}^{k+1} \end{bmatrix} = \begin{bmatrix} d_1 \\ d_2 \\ d_3 \\ \vdots \\ d_{n_x-1} \\ d_{n_x} \end{bmatrix} \quad (3)$$

where the coefficients of the interior nodes ($i = 2, 3, \dots, N - 1$) are

$$a_i = 1/\Delta t + \alpha/\Delta x^2 = 1/\Delta t - (b_i + c_i),$$

$$b_i = c_i = -\alpha/(2\Delta x^2),$$

$$d_i = -c_i u_{i-1}^k + (1/\Delta t + b_i + c_i) u_i^k - b_i u_{i+1}^k.$$

demoCN Code

```
% --- Coefficients of the tridiagonal system
b = (-alfa/2/dx^2)*ones(nx,1); % Super diagonal: coefficients of u(i+1)
c = b; % Subdiagonal: coefficients of u(i-1)
a = (1/dt)*ones(nx,1) - (b+c); % Main Diagonal: coefficients of u(i)
at = (1/dt + b + c); % Coefficient of u_i^k on RHS
a(1) = 1; b(1) = 0; % Fix coefficients of boundary nodes
a(end) = 1; c(end) = 0;
[e,f] = tridiagLU(a,b,c); % Save LU factorization

% --- Assign IC and save BC values in ub. IC creates u vector
x = linspace(0,L,nx)'; u = sin(pi*x/L); ub = [0 0];

% --- Loop over time steps
for k=2:nt
    % --- Update RHS for all equations, including those on boundary
    d = - [0; c(2:end-1).*u(1:end-2); 0] ...
        + [ub(1); at(2:end-1).*u(2:end-1); ub(2)] ...
        - [0; b(2:end-1).*u(3:end); 0];
    u = tridiagLUsolve(e,f,b,d); % Solve the system
end
```


Convergence of FTCS, BTCS and CN

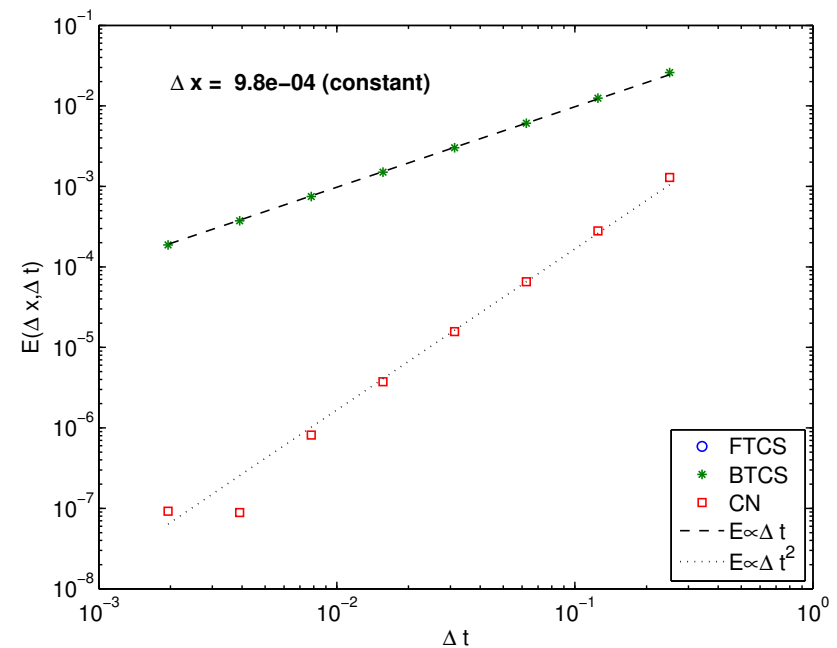
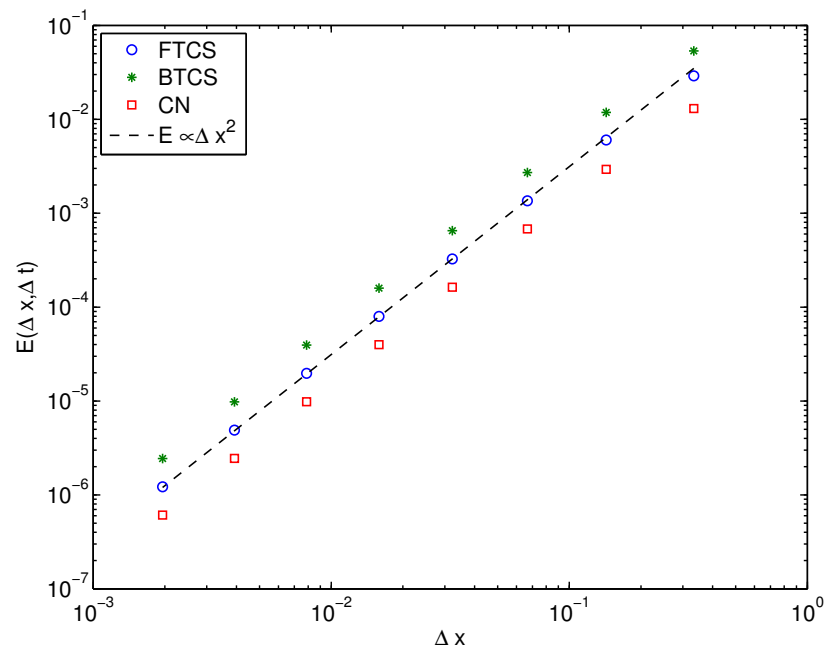
Reduce both dx and dt within the FTCS stability limit

Truncation Errors			Errors				
Scheme	Spatial	Temporal	nx	nt	FTCS	BTCS	CN
FTCS	Δx^2	Δt	4	5	2.903e-02	5.346e-02	1.304e-02
			8	21	6.028e-03	1.186e-02	2.929e-03
BTCS	Δx^2	Δt	16	92	1.356e-03	2.716e-03	6.804e-04
			32	386	3.262e-04	6.522e-04	1.630e-04
C-N	Δx^2	Δt^2	64	1589	7.972e-05	1.594e-04	3.984e-05
			128	6453	1.970e-05	3.939e-05	9.847e-06
			256	26012	4.895e-06	9.790e-06	2.448e-06
			512	104452	1.220e-06	2.440e-06	6.101e-07

Reduce dt while holding dx = 9.775171e-04 (L=1.0, nx=1024) constant

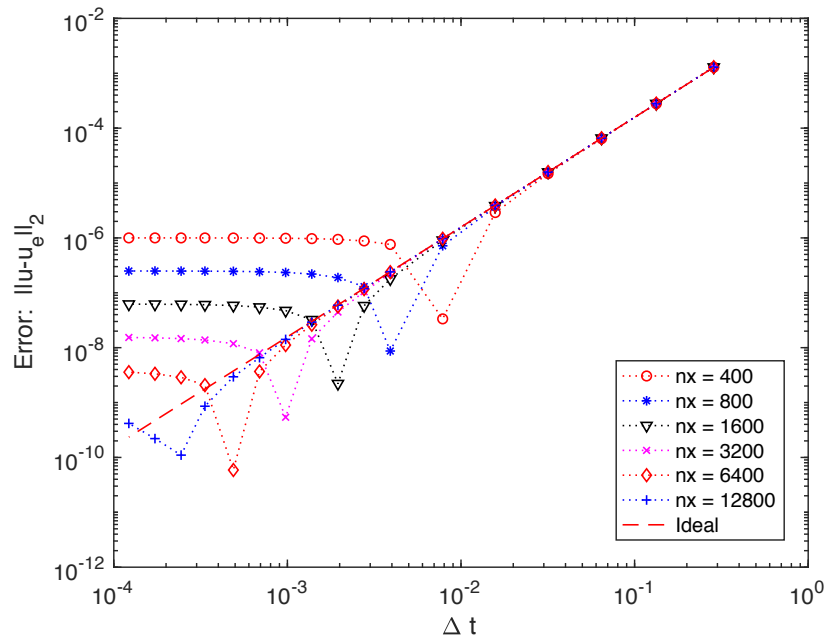
		Errors		
nx	nt	FTCS	BTCS	CN
1024	8	NaN	2.601e-02	1.291e-03
1024	16	NaN	1.246e-02	2.798e-04
1024	32	NaN	6.102e-03	6.534e-05
1024	64	NaN	3.020e-03	1.570e-05
1024	128	NaN	1.502e-03	3.749e-06
1024	256	NaN	7.492e-04	8.154e-07
1024	512	NaN	3.742e-04	8.868e-08
1024	1024	NaN	1.871e-04	9.218e-08

Convergence of FTCS, BTCS and CN



In the plot of truncation error versus Δt (right hand plot), there is an irregularity at $\Delta t \sim 3.9 \times 10^{-3}$. At that level of Δt , and for the chosen Δx (which is constant), the truncation error due to Δx is no longer negligible. Further reductions in Δt alone will not reduce the total truncation error.

Truncation Error is Additive



The truncation error for the finite-difference schemes that we have explored so far are of the form

$$\bar{e} = \mathcal{O}(\Delta t^p) + \mathcal{O}(\Delta x^q)$$

For example, for the Crank-Nicolson scheme, $p = q = 2$.

If Δt is reduced while Δx is held constant, the measured error is reduced until the point that the temporal truncation error is less than the spatial truncation error.

Summary for the Crank-Nicolson Scheme

- The Crank-Nicolson method is more accurate than FTCS or BTCS. Although all three methods have the same spatial truncation error (Δx^2), the better *temporal* truncation error for the Crank-Nicolson method is a big advantage.
- Like BTCS, the Crank-Nicolson scheme is unconditionally stable for the heat equation.
- Like BTCS, a system of equations for the unknown u_i^k must be solved at each time step. The tridiagonal solver for the 1D heat equation obtains an efficient solution of the system of equations.
- The Crank-Nicolson scheme is recommended over FTCS and BTCS.