# Alternative Boundary Condition Implementations for Crank Nicolson Solution to the Heat Equation 

ME 448/548 Notes

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## Overview

1. Goal is to allow Dirichlet, Neumann and mixed boundary conditions
2. Use ghost node formulation

- Preserve spatial accuracy of $\mathcal{O}\left(\Delta x^{2}\right)$
- Preserve tridiagonal structure to the coefficient matrix

3. Implement in a code that uses the Crank-Nicolson scheme.
4. Demonstrate the technique on sample problems

## Heat Transfer Boundary Conditions



1. Prescribe $T_{w}$, a know wall temperature. Maybe $T_{w}=f(t)$.
2. Solve internal $T(x, t)$ field
3. Compute the wall heat flux, $q_{w}$.

4. Prescribe $q_{w}$, a know wall heat flux. Maybe $q_{w}(t)=f(t)$.
5. Solve internal $T(x, t)$ field
6. Compute the wall temperature, $T_{w}$.

7. Prescribe $T_{\infty}$ and $h$. Maybe $T_{w}(t)=f(t)$ and $h=f(t)$.
8. Solve internal $T(x, t)$ field
9. Compute the wall heat flux, $q_{w}$ and wall temperature, $T_{w}$.

## Convective Boundary Condition

The general form of a convective boundary condition is

$$
\begin{equation*}
\left.\frac{\partial u}{\partial x}\right|_{x=0}=g_{0}+h_{0} u \tag{1}
\end{equation*}
$$

This is also known as a Robin boundary condition or a boundary condition of the third kind.

The simplistic implementation is to replace the derivative in Equation (1) with a one-sided difference

$$
\begin{equation*}
\frac{u_{2}^{k+1}-u_{1}^{k+1}}{\Delta x}=g_{0}+h_{0} u_{1}^{k+1} \tag{2}
\end{equation*}
$$

Don't do that! The one-sided difference approximation has a spatial accuracy of $\mathcal{O}(\Delta x)$.

## Introduce a Ghost Node

Imagine that there is a node $\hat{u}_{0}$ that is outside of the domain

this node is used to enforce the boundary condition from Equation (1).
The value $\hat{u}_{0}$ does not explicitly appear in the numerical scheme. We introduce it as a device to introduce a higher order approximation to the gradient at the boundary. It turns out that with algebra, $\hat{u}_{0}$ disappears from the final formulation.

## Use the $\mathbf{B C}$ to compute $\hat{u}_{0}$ by extrapolation

Use a central difference approximation at $x=0$ ( $x=x_{1}$ ) to impose the boundary condition.

$$
\begin{equation*}
\frac{u_{2}-\hat{u}_{0}}{2 \Delta x}=g_{0}+h_{0} u_{1} . \tag{3}
\end{equation*}
$$

The value of $\hat{u}_{0}$ consistent with the boundary condition is

$$
\begin{equation*}
\hat{u}_{0}=u_{2}-2 \Delta x\left(g_{0}+h_{0} u_{1}\right) . \tag{4}
\end{equation*}
$$

Equation (4) allows us to eliminate $\hat{u}_{0}$ at the boundary.

## Equation for $u_{1}$

Evaluate the finite difference form of the heat equation at $x=x_{1}$.

$$
\frac{u_{1}^{k+1}-u_{1}^{k}}{\Delta t}=\theta \alpha\left[\frac{\hat{u}_{0}^{k+1}-2 u_{1}^{k+1}+u_{2}^{k+1}}{\Delta x^{2}}\right]+(1-\theta) \alpha\left[\frac{\hat{u}_{0}^{k}-2 u_{1}^{k}+u_{2}^{k}}{\Delta x^{2}}\right]
$$

Choose $\theta=1 / 2$ and use the formulas for $\hat{u}_{0}$ at time step $k$ and time step $k+1$

$$
\begin{aligned}
\frac{u_{1}^{k+1}-u_{1}^{k}}{\Delta t}=\frac{\theta \alpha}{\Delta x^{2}} & {\left[\begin{array}{|}
u_{2}^{k+1}-2 \Delta x\left(g_{0}^{k+1}+h_{0}^{k+1} u_{1}^{k+1}\right) & \left.-2 u_{1}^{k+1}+u_{2}^{k+1}\right] \\
& +\frac{(1-\theta) \alpha}{\Delta x^{2}}\left[\boxed{u_{2}^{k}-2 \Delta x\left(g_{0}^{k}+h_{0}^{k} u_{1}^{k}\right)}-2 u_{1}^{k}+u_{2}^{k}\right]
\end{array}\right.}
\end{aligned}
$$

The terms in boxes are from the boundary condition

## Rearrange the Equation for $u_{1}$

Algebraically rearranging the preceding equation gives

$$
\begin{equation*}
a_{1} u_{1}^{k+1}+b_{1} u_{2}^{k+1}=d_{1} \tag{5}
\end{equation*}
$$

where

$$
\begin{align*}
a_{1}= & \frac{1}{\Delta t}+\frac{2 \theta \alpha}{\Delta x^{2}}\left(1+\Delta x h_{0}^{k+1}\right)  \tag{6}\\
b_{1}= & -\frac{2 \theta \alpha}{\Delta x^{2}}  \tag{7}\\
d_{1}= & {\left[\frac{1}{\Delta t}-\frac{2(1-\theta) \alpha}{\Delta x^{2}}\left(1+\Delta x h_{0}^{k}\right)\right] u_{1}^{k} }  \tag{8}\\
& \quad+\frac{2(1-\theta) \alpha}{\Delta x^{2}} u_{2}^{k}-\frac{2 \alpha}{\Delta x}\left[\theta g_{0}^{k+1}+(1-\theta) g_{0}^{k}\right]
\end{align*}
$$

These equations define the terms for the first row in the system of equations

## Data structure for implementing alternative BC in the MATLAB code

Store the data defining the boundary condition for both boundaries in a $2 \times 3$ matrix.

The first row has data for $x=0$
The second row has data for $x=L$.

$$
\mathrm{u}^{\prime} \mathrm{boc}=\begin{array}{|l|l|l|}
\hline \text { Type } & \text { Value 1 } & \text { Value } 2 \\
\hline \text { Type } & \text { Value 1 } & \text { Value } 2 \\
\hline & x=0 \\
x=L
\end{array}
$$

Type is a flag with the boundary condition type.

$$
\begin{aligned}
& \text { if } \mathrm{ubc}(\mathrm{~b}, 1)=1 \text {, then } \\
& \qquad \begin{array}{l}
u\left(x_{b}\right)=\text { value } \\
\mathrm{ubc}(\mathrm{~b}, 2)=\text { value of } u \text { at boundary } \\
\mathrm{ubc}(\mathrm{~b}, 3)=\text { not used } \\
\text { if } \mathrm{ubc}(\mathrm{~b}, 1)=2 \text {, then } \\
\partial u /\left.\partial x\right|_{x_{b}}=g+h u\left(x_{b}\right) \\
\operatorname{ubc}(\mathrm{b}, 2)=g \\
\operatorname{ubc}(\mathrm{~b}, 3)=h
\end{array}
\end{aligned}
$$

## Verification: Solve the toy problem on half of the domain

The toy problem used to test the codes

$$
\begin{aligned}
& \frac{\partial u}{\partial t}=\alpha \frac{\partial^{2} u}{\partial x^{2}} \quad t>0, \quad 0 \leq x \leq L \\
& u(0, t)=u(L, t)=0 \\
& u(x, 0)=\sin (\pi x / L)
\end{aligned}
$$

only needs to be solved on one half of the domain

$$
\begin{array}{ll}
\frac{\partial u}{\partial t}=\alpha \frac{\partial^{2} u}{\partial x^{2}} \quad t>0,0 \leq x \leq L / 2 & \frac{\partial u}{\partial t}=\alpha \frac{\partial^{2} u}{\partial x^{2}} \quad t>0, L / 2 \leq x \leq L \\
u(0, t)=0 ;\left.\quad \frac{\partial u}{\partial x}\right|_{L / 2}=0 & \left.\frac{\partial u}{\partial x}\right|_{L / 2}=0 \quad u(L, t)=0 \\
u(x, 0)=\sin (\pi x / L) & u(x, 0)=\sin (\pi x / L)
\end{array}
$$

## Verification: Solve the toy problem on half of the domain

Use ubc matrix to specify boundary conditions.
For the half-problem on $0 \leq x \leq L / 2$ :

$$
\begin{aligned}
& u(0, t)=0 \quad \Longrightarrow \quad \begin{array}{l}
\operatorname{ubc}(1,1)=1, \quad \text { boundary type } \\
\operatorname{ubc}(1,2)=0, \quad \text { value of } u \text { at } x=0 \\
\operatorname{ubc}(1,3)=0, \quad \text { not used }
\end{array} \\
& \left.\frac{\partial u}{\partial x}\right|_{L / 2}=0 \quad \Longrightarrow \quad \begin{array}{l}
\operatorname{ubc}(2,1)=2, \quad \text { boundary type } \\
\operatorname{ubc}(2,2)=0, \quad \text { value of } g_{L / 2} \\
\operatorname{ubc}(2,3)=0, \quad \text { value of } h_{L / 2}
\end{array}
\end{aligned}
$$

## Verification: Solve the toy problem on half of the domain

For the half-problem on $L / 2 \leq x \leq L$ :

$$
\begin{aligned}
& \left.\frac{\partial u}{\partial x}\right|_{L / 2}=0 \quad \Longrightarrow \quad \begin{array}{l}
\operatorname{ubc}(1,1)=2, \\
\operatorname{ubc}(1,2)=0, \\
\operatorname{ubc}(1,3)=0,
\end{array} \\
& \text { boundary type of } g_{0} \\
& u(L, t)=0 \quad \Longrightarrow \quad \begin{array}{l}
\text { value of } h_{0}
\end{array} \\
& \operatorname{ubc}(2,1)=1, \quad \text { boundary type } \\
& \operatorname{ubc}(2,2)=0, \quad \text { value of } u \text { at } x=L \\
& \operatorname{ubc}(2,3)=0, \quad \text { not used }
\end{aligned}
$$

## Verification: Solve the toy problem on half of the domain

Output of demoCNBC


## Solve the hot pot problem



Contact resistance at $x=0$

$$
\begin{gathered}
-\left.k \frac{\partial T}{\partial x}\right|_{x=0}=q_{t}=h_{t}\left(T_{p}-T_{0}\right) \\
\left.\Longrightarrow \frac{\partial T}{\partial x}\right|_{x=0}=-\frac{h_{t} T_{p}}{k}+\frac{h_{t}}{k} T_{0}
\end{gathered}
$$

Matlab boundary matrix for $x=0$
$\operatorname{ubc}(1,1)=2, \quad$ boundary type $\operatorname{ubc}(1,2)=-h_{t} T_{p} / k, \quad$ value of $g_{0}$ $\operatorname{ubc}(1,3)=h_{t} / k$, value of $h_{0}$

## Solve the hot pot problem



Convective resistance at $x=L$

$$
\begin{gathered}
-\left.k \frac{\partial T}{\partial x}\right|_{x=L}=q_{b}=h_{b}\left(T_{L}-T_{\mathrm{air}}\right) \\
\left.\Longrightarrow \frac{\partial T}{\partial x}\right|_{x=L}=\frac{h_{b} T_{\mathrm{air}}}{k}-\frac{h_{b}}{k} T_{L}
\end{gathered}
$$

Matlab boundary matrix for $x=L$
$\operatorname{ubc}(2,1)=2$,
$\operatorname{ubc}(2,2)=h_{b} T_{\text {air }} / k$,
boundary type value of $g_{L}$
$\operatorname{ubc}(2,3)=-h_{b} / k$, value of $h_{L}$

## Solve the hot pot problem

Core of demoHotPot.m

```
% --- Define physical properties for table and boundary conditions
rho = 545; % density of oak (kg/m^3)
k = 0.17; % thermal conductivity of oak, across the grain (W/m/C)
c = 2385; % specific heat capacity of oak (J/kg/K)
alfa = k/rho/c; % thermal diffusivity (m^2/s)
L = 2e-2; % Table thickness (m)
% --- Use relaxation time of table material to specify time step size
tau = L^2/alfa; % Relaxation time for the heat condution (s)
dt = tau/1000; % Time step (s)
nt = ceil(tmax/dt); % Number of time steps
% --- Specify initial and boundary conditions
u0 = Tair*ones(nx,1);
ubc = [2 (-htop*Tp/k) (htop/k); 2 (hbot*Tair/k) (-hbot/k)];
% --- Solve the heat equation and plot the results
[U,x,t] = heatCNBC(nx,nt,ubc,u0,L,tmax,alfa);
plotHeat(U,100*x,t,floor(nt/5))
xlabel('x (cm)'); ylabel('T ({{}^\circ}C)'); ylim([Tair-5, Tp])
```


## Solve the hot pot problem

>> demoHotPot


## Solve the hot pot problem

>> demoHotPot(1200)


## Solve the hot pot problem



