

Introduction to Quantum logic (2)

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Classical Logic Circuits

Circuit behavior is governed implicitly

by classical physics

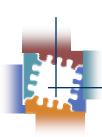
Signal states are simple <u>bit vectors</u>,

e.g. X = 01010111

Operations are defined by <u>Boolean Algebra</u>







Quantum Logic Circuits

- Circuit behavior is governed explicitly
 by quantum mechanics
- Signal states are vectors interpreted as a superposition of binary "qubit" vectors with complex-number coefficients

$$|\Psi\rangle = \sum_{i=0}^{2^{n}-1} c_{i} |i_{n-1}i_{n-1}...i_{0}\rangle$$

Operations are defined by <u>linear algebra over</u> <u>Hilbert Space</u> and can be represented by <u>unitary</u> <u>matrices with complex elements</u>







Signal state (one qubit)

$$\ket{0}$$
 Corresponsd to $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$

$$\begin{vmatrix} 1 \end{vmatrix}$$
 Corresponsd to $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$

$$\alpha_0\big|0\big\rangle+\alpha_1\big|1\big\rangle \qquad \text{Corresponsd to} \qquad \alpha_0\bigg(1\\0\bigg)+\alpha_1\bigg(1\\0\bigg)=\bigg(\alpha_0\\\alpha_1\bigg)\\ \big|\alpha_0\big|^2+\big|\alpha_1\big|^2=1$$







More than one qubit

If we concatenate two qubits

$$\left(\alpha_{0}\left|0\right\rangle+\alpha_{1}\left|1\right\rangle\right)\left(\beta_{0}\left|0\right\rangle+\beta_{1}\left|1\right\rangle\right) \qquad \qquad \frac{\left|\alpha_{0}\right|^{2}+\left|\alpha_{1}\right|^{2}=1}{\left|\beta_{0}\right|^{2}+\left|\beta_{1}\right|^{2}=1}$$

We have a 2-qubit system with 4 basis states

$$|0\rangle|0\rangle = |00\rangle \qquad |0\rangle|1\rangle = |01\rangle \qquad |1\rangle|0\rangle = |10\rangle \qquad |1\rangle|1\rangle = |11\rangle$$

And we can also describe the state as

$$\alpha_0 \beta_0 |00\rangle + \alpha_0 \beta_1 |01\rangle + \alpha_1 \beta_0 |10\rangle + \alpha_1 \beta_1 |11\rangle$$

Or by the vector
$$(\alpha_0 \beta_0 | 00) + \alpha_0 \beta_1 | 01) + \alpha_1 \beta_0 | 10) + \alpha_1 \beta_1 | 11)$$

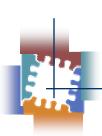
$$(\alpha_0 \beta_0 | \alpha_0 \beta_1 | \alpha_0 \beta_1 | \alpha_0 \beta_0 | \alpha_0 \beta_0 | \alpha_0 \beta_1 | \alpha_0 \beta_0 | \alpha_0 | \alpha_0 \beta_0 | \alpha_0 | \alpha_0$$



Tensor product







Quantum Operations

Any linear operation that takes states

$$|\alpha_0|0\rangle + |\alpha_1|1\rangle$$
 satisfying $|\alpha_0|^2 + |\alpha_1|^2 = 1$

and maps them to states

$$|\beta_0|0\rangle + |\beta_1|1\rangle$$
 satisfying $|\beta_0|^2 + |\beta_1|^2 = 1$

must be UNITARY

$$\overline{U}^{t}U = I$$

$$\Leftrightarrow \overline{U}^t = U^{-1}$$







Find the quantum gate(operation)

$$\overline{U}^{t}U = I$$

$$\Leftrightarrow \overline{U}^{t} = U^{-1}$$

From upper statement

We now know the necessities

- 1. A matrix must has inverse, that is reversible.
- 2. Inverse matrix is the same as \overline{U}^t .

So, reversible matrix is good candidate for quantum gate



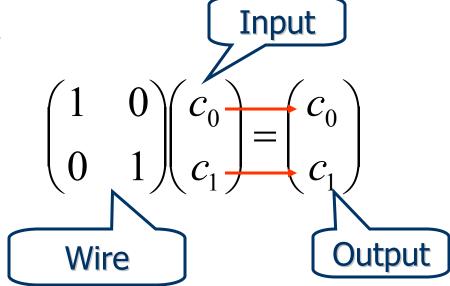




One-Input gate: Wire



By matrix





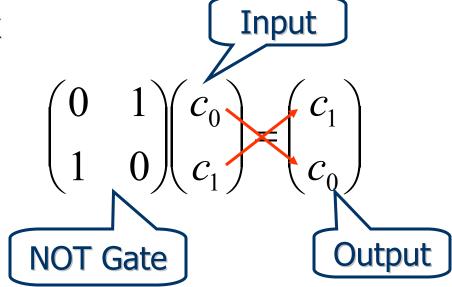




One-Input gate: NOT



By matrix

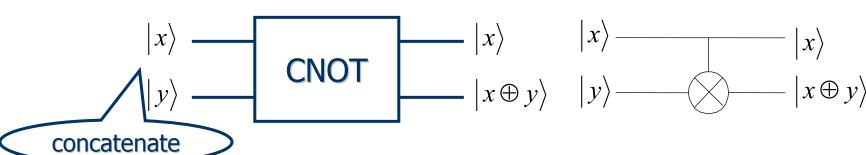






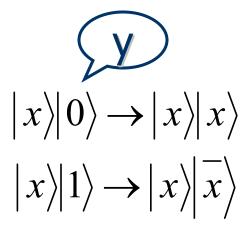


Two-Input Gate: Controlled NOT (Feynman gate)



By matrix

$$\begin{pmatrix}
1 & 0 & 0 & 0 \\
0 & 1 & 0 & 0 \\
0 & 0 & 0 & 1 \\
0 & 0 & 1 & 0
\end{pmatrix}
\begin{pmatrix}
00 \\
01 \\
10 \\
11
\end{pmatrix}
=
\begin{pmatrix}
00 \\
01 \\
11 \\
10
\end{pmatrix}$$

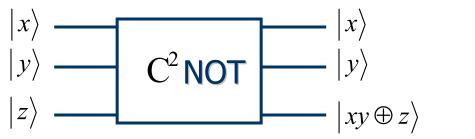








3-Input gate: Controlled CNOT (Toffoli gate)



$$\begin{vmatrix} x \rangle & & |x \rangle \\ |y \rangle & & |y \rangle \\ |z \rangle & & |xy \oplus z \rangle \end{vmatrix}$$

❖ By matrix

$$\begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & 1 & 0 \\ \end{pmatrix}$$

$$\begin{pmatrix}
0 & 0 & 0 \\
0 & 0 & 1 \\
0 & 1 & 0 \\
0 & 1 & 1 \\
1 & 0 & 0 \\
1 & 0 & 1 \\
1 & 1 & 0 \\
1 & 1 & 1
\end{pmatrix}$$

$$\begin{pmatrix}
0 & 0 & 0 \\
0 & 0 & 1 \\
0 & 1 & 0 \\
0 & 1 & 1 \\
1 & 0 & 0 \\
1 & 0 & 1 \\
1 & 1 & 1 \\
1 & 1 & 0
\end{pmatrix}$$





Quantum circuit

